THE EXponential STABILITY OF A COUPLED
HYPERBOLIC/PARABOLIC SYSTEM ARISING IN STRUCTURAL
ACOUSTICS

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The Exponential Stability of a Coupled Hyperbolic/Parabolic System Arising in Structural Acoustics

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Abstract

We show here the uniform stabilization of a coupled system of hyperbolic and parabolic PDE's which describes a particular fluid/structure interaction system. This system has the wave equation, which is satisfied on the interior of a bounded domain $\Omega$, coupled to a "parabolic-like" beam equation holding on $\partial\Omega$, and wherein the coupling is accomplished through velocity terms on the boundary. Our result is an analog of that in [12] which shows the exponential stability of the wave equation via Neumann feedback control, and like that work, depends upon a trace regularity estimate for solutions of hyperbolic equations.

1 Introduction

1.1 Statement of the Problem and Motivation

Let $\Omega$ be a bounded domain of $\mathbb{R}^n$, $n \geq 2$, with Lipshitz boundary $\Gamma = \Gamma_0 \cup \Gamma_1$, $\Gamma_i$ open, $\Gamma_0$ and $\Gamma_1$ both nonempty, with $\Gamma_0 \cap \Gamma_1 = \emptyset$. This paper is a continuation of our study, initiated in ([1]), of the following system consisting of a coupling between a wave and plate-like equation:

\[
\begin{align*}
z_{tt} &= \Delta z & \text{on } (0, \infty) \times \Omega \\
z(0, x) &= z_0, \quad z_t(0, x) = z_1 & \text{on } \Omega \\
z(t, x) &= 0 & \text{on } (0, \infty) \times \Gamma_0 \\
\frac{\partial z(t, x)}{\partial \nu} + \alpha z_t(t, x) &= v_t & \text{on } (0, \infty) \times \Gamma_1 \text{ with } \alpha \geq 0; \\
v_{tt} &= -\Delta^2 v - \Delta^2 v_t - z_t & \text{on } (0, \infty) \times \Gamma_1 \\
v(0, x) &= v_0, \quad v_t(0, x) = v_1 & \text{on } \Gamma_1 \\
v(t, x) &= \frac{\partial v(t, x)}{\partial \nu} = 0 & \text{on } (0, \infty) \times \partial \Gamma_1.
\end{align*}
\]

(1)

(2)

Note how this coupling above of two qualitatively different equations is accomplished by the velocity terms $z_t$ and $v_t$ on the active portion of the boundary $\Gamma_1$.

In [2], issues of well-posedness for (1)-(2) were settled, with initial data $[z_0, v_0] \equiv [z_0, z_1, v_0, v_1]$ determining the solution $[z, v] \equiv [z, z_t, v, v_t]$ to be in $H^{1}_{\Gamma_0}(\Omega) \times L^2(\Omega) \times H^2_0(\Gamma_1) \times L^2(\Gamma_1)$ (where $H^{1}_{\Gamma_0}(\Omega) = \{ z \in H^1(\Omega) \, \| \, z = 0 \text{ on } \Gamma_0 \}$. Here, we are concerned with the exponential decay of the

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solution \([\mathbf{x}, \mathbf{v}]\) to (1)-(2). Specifically, we wish to know: Defining the energy \(E(\mathbf{x}, \mathbf{v}, t)\) of the system as

\[
E(\mathbf{x}, \mathbf{v}, t) = \int_{\Omega} \left( |\nabla z(t)|^2 + |z_t(t)|^2 \right) d\Omega + \int_{\Gamma_1} \left( |\Delta v(t)|^2 + |v_t(t)|^2 \right) d\Gamma_1,
\]

(3)
does there exist positive constants \(C\) and \(\omega\) such that

\[
E(\mathbf{x}, \mathbf{v}, t) \leq Ce^{-\omega t} \left\| \begin{bmatrix} z_0 \\ \frac{\partial z_0}{\partial \nu} \end{bmatrix} \right\|^2
\]

(where the norm above denotes that of the \(H^1_0(\Omega) \times L^2(\Omega) \times H^2_0(\Gamma_1) \times L^2(\Gamma_1)\)-topology). This "structural acoustics" model is a variation of that derived by H.T. Banks et al (see [4],[5]) to mathematically describe the interaction between an acoustic field and its vibrating boundary, a physical phenomenon much studied nowadays within the realm of smart materials and structures and its accompanying numerical analysis. A simple PDE argument will reveal that the original system of Banks et al does not exhibit uniform decay, and hence the necessity for the supplanting with (1)-(2). We furthermore focus here on the case that the parameter \(\alpha > 0\).

A demonstration of exponential stability for the system (1)-(2) has important physical implications, as one would consequently be free to study the associated Linear Quadratic Regulator Problem (LQR) on Infinite Horizon. In the LQR for the structural acoustics model, boundary control is implemented via the placement of linear combinations of derivatives of delta functions in the beam equation of (2), so as to model the use of piezoelectric ceramic patches in inducing acoustic noise reduction (the LQR for finite time has been given a thorough treatment in [1]). Note here that the input operator which models the control action is "badly" unbounded, and the LQR is consequently not amenable to the recently developed treatments in [9]. The exponential stability of (1)-(2) is requisite in the analysis of the LQR for infinite time, which in turn could potentially yield a viable numerical approach (via a formulation of the appropriate Algebraic Riccati Equation) for obtaining approximations of the structural acoustics control problem. Moreover, the uniform stabilization of (1)-(2) is indispensable in future considerations of non-linear versions of the model.

1.2 Preliminaries

In dealing with (1)-(2), we will consider throughout its equivalence with an abstract evolution equation, for the defining of which we will need the following background material:

- Let the operator \(A : L^2(\Omega) \supset D(A) \to L^2(\Omega)\) be defined by

\[
Az = -\Delta z, \quad D(A) = \left\{ z \in H^2(\Omega) \; \exists \; z|_{\Gamma_0} = 0, \left. \frac{\partial z}{\partial \nu} \right|_{\Gamma_1} = 0 \right\}.
\]

(4)

Note that \(A\) is self-adjoint, positive definite, and hence the fractional powers of \(A\) are well defined.

- By [10], we have the following characterization:

\[
D(A^{\frac{1}{2}}) = H^1_0(\Omega) = \left\{ z \in H^1(\Omega) \; \exists \; z = 0 \; {\text{on}} \; \Gamma_0 \right\},
\]

(5)

with \(|z|^2_{D(A^{\frac{1}{2}})} = \left\| A^{\frac{1}{2}} z \right\|_{L^2(\Omega)}^2 = \int_{\Omega} |\nabla z|^2 d\Omega = \|z\|_{H^1_0}^2 \; \forall \; z \in D(A^{\frac{1}{2}}),\]

(6)

where the last equality in (6) follows from Poincaré's inequality.
• We define the map $N$ by

$$
\begin{aligned}
z = Ng & \iff \\
\begin{cases}
\Delta z = 0 & \text{on } \Omega \\
z|_{\Gamma_0} = 0 & \text{on } \Gamma_0 \\
\frac{\partial z}{\partial \nu}_{|_{\Gamma_1}} = g & \text{on } \Gamma_1;
\end{cases}
\end{aligned}
$$

(7)

elliptic theory will then yield that

$$
N \in L(L^2(\Gamma_1), D(A^\frac{1}{2} - \epsilon)) \quad \forall \epsilon > 0.
$$

(8)

• Let $\gamma : H^1(\Omega) \to H^\frac{1}{2}(\Gamma_1)$ be the restriction to $\Gamma_1$ of the familiar Sobolev trace map; viz.

$$
\forall z \in H^1(\Omega), \quad \gamma(z) = \begin{cases}
z|_{\Gamma_1} & \text{on } \Gamma_1 \\
0 & \text{on } \Gamma_0.
\end{cases}
$$

(9)

Then as is shown in [14], we have

$$
N^*A = \gamma(z) \quad \forall z \in D(A^\frac{1}{2}).
$$

(10)

• We set $\hat{A} : L^2(\Gamma_1) \supset D(\hat{A}) \to L^2(\Gamma_1)$ to be

$$
\hat{A} = \Delta^2, \quad D(\hat{A}) = H^4(\Gamma_1) \cap H^2_0(\Gamma_1);
$$

(11)

$\hat{A}$ is also self–adjoint, positive definite, and again by [10], we have the characterization

$$
D(\hat{A}^\theta) = H^{4\theta}_0(\Gamma_1), \quad 0 < \theta < 1,
$$

(12)

with

$$
\|\hat{A}^{\frac{1}{2}} v\|_{L^2(\Gamma_1)}^2 = \int_{\Gamma_1} |\Delta v|^2 \, d\Omega = \|v\|_{H^2(\Gamma_1)}^2 \quad \forall v \in D(\hat{A}^{\frac{1}{2}}).
$$

(13)

• We define the energy spaces

$$
H_1 \equiv D(A^\frac{1}{2}) \times L^2(\Omega); \quad H_0 \equiv D(\hat{A}^{\frac{1}{2}}) \times L^2(\Gamma_1).
$$

(14)

(15)

• We define $A_1 : H_1 \supset D(A_1) \to H_1$ and $A_0 : H_0 \supset D(A_0) \to H_0$ to be

$$
A_1 \equiv \begin{bmatrix} 0 & I \\ -A & -\alpha AN^* A \end{bmatrix} \quad \text{with}
$$

$$
D(A_1) = \left\{ [z_1, z_2]^T \in \left[ D(A^\frac{1}{2}) \right]^2 \ | z_1 + \alpha NN^* A z_2 \in D(A) \right\};
$$

(17)

$$
A_0 \equiv \begin{bmatrix} 0 & I \\ -\hat{A} & -\hat{A} \end{bmatrix} \quad \text{with}
$$

$$
D(A_0) = \left\{ [v_1, v_2]^T \in \left[ D(\hat{A}^\frac{1}{2}) \right]^2 \ | v_1 + v_2 \in D(\hat{A}) \right\}.
$$

(19)
With $\gamma$ as defined in (28) we set

$$C = \begin{bmatrix} 0 & 0 \\ 0 & \gamma^* \end{bmatrix} \in \mathcal{L}(H_0, D(A^{1/2}) \times D(A^{1/2})') \tag{20}.$$ 

With the above operator definitions, we set

$$A = \begin{bmatrix} A_1 & C \\ -C^* & A_0 \end{bmatrix} \quad \text{with}$$

$$D(A) = \left\{ (z_1, z_2, v, v_2)^T \in D(A^{1/2})^2 \times D(A^{1/2})^2 \right\}$$

such that $-z_1 - \alpha NN^* A z_2 + N v_2 \in D(A)$ and $v_1 + v_2 \in D(\tilde{A}) \tag{21}$. 

If we take the initial data $[\overline{z}_0, \overline{v}_0]$ to be in $H_1 \times H_0$ (as defined in (14)-(15)), we can use the definitions above to rewrite (1)-(2) abstractly as

$$\frac{d}{dt} \begin{bmatrix} z \\ z_t \\ v \\ v_t \end{bmatrix} = A \begin{bmatrix} \overline{z} \\ \overline{v} \end{bmatrix} \tag{22}$$

$$[z(0), v(0)] = [\overline{z}_0, \overline{v}_0].$$

**Remark 1** The structure of $A$ given in (22) clearly reflects the coupled nature of this particular system; the operator $A_1$ which models hyperbolic dynamics is linked via an unbounded coupling with the “elastic” operator $A_0$ which exhibits parabolic characteristics, and the coupling here is accomplished by the “trace” operators $C$ and $C^*$. 

From (22), the differential equations in (1)-(2) then have the following abstract representation:

$$z_{tt} = -Az - ANN^* Az_t + AN v_t \quad \text{on} \ (0, \infty) \times \Omega; \tag{23}$$

$$v_{tt} = -\tilde{A} v - \tilde{A} v_t - N A^* z_t \quad \text{on} \ (0, \infty) \times \Gamma_1. \tag{24}$$

Regarding the well–posedness and strong stability of (1)-(2) and its equivalent form (22), we have the following recent result:

**Theorem A** (see [2]). With $\alpha \geq 0$ in (1),

(i) $A$ given by (21) generates a $C_0-$semigroup of contractions $\{e^{At}\}_{t \geq 0}$ on the energy space $H_1 \times H_0$.

(ii) The semigroup $\{e^{At}\}_{t \geq 0}$ is strongly stable; that is,

$$\forall \ [\overline{z}_0, \overline{v}_0] \equiv [z_0, z_1, v_0, v_1] \in H_1 \times H_0, \text{ one has} \tag{25}$$

$$\lim_{t \to \infty} e^{At} \begin{bmatrix} \overline{z}_0 \\ \overline{v}_0 \end{bmatrix} \to 0.$$ 

**1.3 Literature**

The exponential stability for the individual components $A_1$ and $A_0$ have been well–established these past few years, but that of the entire structure $A$ has not been addressed. Concerning the beam equation modelled by the “elastic” operator $A_0$, we have the result of S. Chen and Triggiani in [8] that $A_0$ generates an analytic semigroup, which automatically provides for the exponential decay of
the solution \([v, v_t]\) of the second-order system
\[
v_{tt} = -\Delta^2 v - \Delta^2 v_t \text{ on } (0, \infty) \times \Gamma_1
\]
\[
v(t, x) = v(t, x) = \frac{\partial v(t, x)}{\partial n} = \frac{\partial v(t, x)}{\partial \nu} = 0 \text{ on } \partial \Gamma_1
\]
\([v(0), v_t(0)] = [v_0, v_1] \in H_1 \times H_0.
\]
For the wave equation with \(L^2(0, T; L^2(\Omega))\)-Neumann feedback control; viz.
\[
z_{tt} = \Delta z \quad \text{on } (0, \infty) \times \Omega
\]
\[
z(0, x) = z_0, \quad z_t(0, x) = z_1 \quad \text{on } \Omega
\]
\[
z(t, x) = 0 \quad \text{on } (0, \infty) \times \Gamma_0
\]
\[
\frac{\partial z(t, x)}{\partial \nu} = -\alpha z_t(t, x) \quad \text{on } (0, \infty) \times \Gamma_1 \quad \alpha > 0;
\]
(27)
G. Chen in [7] proved the exponential stability of solutions (27) under the geometrical conditions that \(\Omega\) be "star-shaped". Lagnese in [11], and subsequently, Triggiani in [14] through an alternate proof, showed the uniform stabilization of (27) under the lessened constraint that there exist a \(C^2(\overline{\Omega})\)-vector field \(h(x)\) such that
(h.i) \(h \cdot \nu \leq 0\) on \(\Gamma_0\)
where \(\nu\) denotes the unit-normal vector to \(\Gamma\);
(h.ii) \(h\) is parallel to \(\nu\) on \(\Gamma_1\);
(h.iii) The Jacobian matrix \(H(x)\) of \(h(x)\) is uniformly positive definite on \(\overline{\Omega}\).

Also, C. Bardos, G. Lebeau and J. Rauch in [6] have derived stability results for wave equations with more general boundary conditions than those in (27), under the assumptions of geometric optics; however the techniques used in the proofs therein are not easily adaptable to our particular situation, based as they are on microlocal analysis and the propagation of singularities. Instead, we shall use the approach of I. Lasiecka and R. Triggiani in [12], who have shown the exponential decay of solutions of (27) without the constraint (h.ii). This result is proved by using the standard multipliers \(h \cdot \nabla z\) and \(z \div h\), and invoking a deep (pseudodifferential) trace estimate which we state here for future reference:

**Lemma A**(see [12]). Let \(\epsilon > 0\) be arbitrarily small. Let \(z\) solve an arbitrary second-order hyperbolic equation on \((0, T)\) with smooth space-dependent coefficients. Then with \(Q_T = (0, T) \times \Omega\),

\[
\int_0^T \int_{\Gamma_1} \left( \frac{\partial z}{\partial \tau} \right)^2 d\Gamma_1 dt < C_{T, \epsilon} \left\{ \int_0^T \int_{\Gamma_1} \left( \left( \frac{\partial z}{\partial \nu} \right)^2 + z_t^2 \right) d\Gamma_1 dt + \|z\|_{H^{1+\epsilon}(Q_T)}^2 \right\},
\]

(28)
where the \(\epsilon\) on the left of (28) need not be the same as the \(\epsilon\) for the \(Q_T\)-norm on the right, and where \(\frac{\partial}{\partial \tau}\) denotes the tangential, and \(\frac{\partial}{\partial \nu}\) the co-normal derivative.

It is this control of the tangential derivative provided above that allows one to forego the condition (h.ii) and generate the desired bound on the energy. In what follows, we will use critically the fact that **Lemma A** is applicable to the wave equation given in (1).
1.4 Statement of Main Result

Theorem 1 With \( \alpha > 0 \) in (1), suppose there exists a \( C^2(\Omega) \) vector field satisfying (h.i) and (h.iii) only. Then, the semigroup \( \{ e^{At} \}_{t \geq 0} \) generated by the operator \( A \) (defined in (21)) is exponentially stable; that is, there exists positive constants \( C \) and \( \omega \) such that the solution \( [\overline{z}, \overline{v}] \) of (1)–(2) satisfies

\[
E(\overline{z}, \overline{v}, t) = \left\| e^{At} \left[ \begin{array}{c} \overline{z} \\ \overline{v} \end{array} \right] \right\|^2_{H_1 \times H_0} \leq Ce^{-\omega t} \left\| \begin{array}{c} \overline{z}_0 \\ \overline{v}_0 \end{array} \right\|^2_{H_1 \times H_0}. \tag{29}
\]

Note that the proof of Theorem 1.1 is independent of the strong stability result posted in Lemma A, wherein there is no imposition of geometrical conditions. As will be explained below, to demonstrate the exponential stability, it will suffice to show that there exists a \( T, 0 < T < \infty \) and corresponding constant \( C_T \) such that

\[
E(\overline{z}, \overline{v}, T) \leq C_T \int_0^T \left( \|z_t\|^2_{L^2(\Gamma_1)} + \left\| \mathbb{A}^{\frac{1}{2}} v_t \right\|^2_{L^2(\Gamma_1)} \right) dt. \tag{30}
\]

To obtain (30), we will rely on the strong damping provided by the beam equation in (2) combined with a multiplier method for the wave equation in (1) to extract a preliminary upper bound on the energy \( E(\overline{z}, \overline{v}, T) \). This upper bound, besides containing the RHS of (30), also includes the tangential derivative of \( z \) and lower order terms. We then use Lemma A, in a very similar way as was done in [12], to estimate \( \frac{\partial z}{\partial t} \) in terms of the RHS of (30) and more lower order terms, and finally eliminate the lower order terms through a compactness/uniqueness argument.

Remark 2 In estimating the energy contribution of (1), one could also proceed as in [11] and [14] to eventually arrive at the uniform stabilization of (1)–(2) under all the geometrical conditions (h.i)–(h.iii). To reiterate, it is the abstract trace estimate (28) which helps to yield the stronger result by eliminating the condition (h.ii).

In proving Theorem 1, we will, without loss of generality, take \( \alpha \equiv 1 \), as one will see in the proofs below that the value of \( \alpha \) is irrelevant, so long as it is positive.

2 Proof of Main Result

Throughout, the initial data \( [\overline{z}_0, \overline{v}_0] \) is taken to be in \( D(A) \), which provides that \( [z, z_t, v, v_t] \in C([0, T]; D(A)) \), and \( [z_t, z_{tt}, v_t, v_{tt}] \in C([0, T]; H_1 \times H_0) \). Proving the results below and subsequently Theorem 1.1 in this special case will be conclusive, as we can then extend the results by density to hold for all initial data in \( H_1 \times H_0 \).

We first give a preliminary result concerning the \( a \text{ priori} \) regularity of the velocity terms \( z_t \) and \( v_t \) which will be used frequently in the work ahead.

Proposition 1 With \( [\overline{z}, \overline{v}] \) the solution of (1)–(2) (guaranteed by Theorem A.(i)), we have

The map \( \{ \overline{z}_0, \overline{v}_0 \} \rightarrow \{ z_t|_{\Gamma_1}, v_t \} \in \mathcal{L} \left( H_1 \times H_0, L^2 \left( 0, \infty; L^2(\Gamma_1) \times D(\mathbb{A}^{\frac{1}{2}}) \right) \right) \). Indeed, we have \( \forall 0 < T < \infty \),

\[
2 \int_0^T \left( \| z_t|_{\Gamma_1} \|^2_{L^2(\Gamma_1)} + \left\| \mathbb{A}^{\frac{1}{2}} v_t \right\|^2_{L^2(\Gamma_1)} \right) dt = E(\overline{z}, \overline{v}, 0) - E(\overline{z}, \overline{v}, T). \tag{31}
\]
Proof: We have by multiplying (23) by $z_t$, (24) by $v_t$, and integrating from 0 to $T$:

$$\int_0^T \|N^* A z_t\|^2_{L^2(\Gamma_1)} dt = \int_0^T \langle AN v_t, z_t \rangle = [D(A^\frac{1}{2})'] \times D(A^\frac{1}{2})$$

$$+ \frac{1}{2} \left( \| A^\frac{1}{2} z_0 \|^2_{L^2(\Omega)} + \| z_t \|^2_{L^2(\Omega)} - \| A^\frac{1}{2} z(T) \|^2_{L^2(\Omega)} - \| z_t(T) \|^2_{L^2(\Omega)} \right) ;$$

$$\int_0^T \| A^\frac{1}{2} v_t \|^2_{L^2(\Gamma_1)} - t \left( N^* A z_t, v_t \right)_{L^2(\Gamma_1)} dt$$

$$+ \frac{1}{2} \left( \| A^\frac{1}{2} v_0 \|^2_{L^2(\Gamma_1)} + \| v_t \|^2_{L^2(\Gamma_1)} - \| A^\frac{1}{2} v(T) \|^2_{L^2(\Gamma_1)} - \| v_t(T) \|^2_{L^2(\Gamma_1)} \right) .$$

Considering the definition of $E$ given in (3) and the characterizations (6), (13) and (10), the desired relation (31) is obtained after the addition of (32)–(33). The asserted continuity of the map \( \{ x_0, \bar{v}_0 \} \rightarrow \{ z_t |_{\Gamma_1}, v_t \} \) is consequently deduced from (31) and the contraction of the semigroup \( \{ e^{A t} \} \_{t \geq 0} . \)

2.1 Proof of Theorem 1.1

A standard argument has that to show the exponential decay (29), it suffices to show that exists a time $0 < T < \infty$ and a corresponding positive constant $C_T$ such that for all initial data in $H_1 \times H_0$:

$$E(\overline{\mathcal{V}}, \overline{\mathcal{V}}, T) \leq \eta E(\overline{\mathcal{V}}, \overline{\mathcal{V}}, 0) \quad \forall \eta < 1; \quad (34)$$

given Proposition 2.1, it will in turn suffice to show that there exists a time $0 < T < \infty$ and a corresponding positive constant $C_T$ such that

$$E(\overline{\mathcal{V}}, \overline{\mathcal{V}}, T) \leq C_T \int_0^T \left[ \| z_t |_{\Gamma_1} \|^2_{L^2(\Gamma_1)} + \| A^\frac{1}{2} v_t \|^2_{L^2(\Gamma_1)} \right] dt, \quad (35)$$

to which end we proceed to work.

We will make use throughout of the denotations $Q_T := (0, T) \times \Omega$, $\Sigma_T := (0, T) \times \Gamma$ and $\Sigma_iT := (0, T) \times \Gamma_i, i = 0, 1.$

Lemma 1 There exists a positive constant $C$, independent of time, such that $\forall 0 < T < \infty$

$$\int_0^T \left[ \| A^\frac{1}{2} v \|^2_{L^2(\Gamma_1)} + \| v_t \|^2_{L^2(\Gamma_1)} \right] dt \leq C \left[ E(\overline{\mathcal{V}}, \overline{\mathcal{V}}, T) + E(\overline{\mathcal{V}}, \overline{\mathcal{V}}, 0) + \int_0^T \left[ \| z_t |_{\Gamma_1} \|^2_{L^2(\Gamma_1)} + \| A^\frac{1}{2} v_t \|^2_{L^2(\Gamma_1)} \right] dt + \| v \|^2_{L^2(\Gamma_1)} \right] . \quad (36)$$

Proof: Trivially, we have that

$$\int_0^T \| v_t \|^2_{L^2(\Gamma_1)} dt \leq \int_0^T \| A^\frac{1}{2} v_t \|^2_{L^2(\Gamma_1)} dt. \quad (37)$$

Moreover, multiplying (24) by $v$ and integrating from 0 to $T$ yields

$$\int_0^T \| A^\frac{1}{2} v \|^2_{L^2(\Gamma_1)} dt = \int_0^T \| v_t \|^2_{L^2(\Gamma_1)} dt - \left[ (v_t, v)_{L^2(\Gamma_1)} \right]_0^T$$

$$+ \frac{1}{2} \left[ \| A^\frac{1}{2} v \|^2_{L^2(\Gamma_1)} \right]_0^T - \int_0^T (N^* A z_t, v)_{L^2(\Gamma_1)} dt ; \quad (38)$$
using (37). Poincare’s inequality, the definition of the energy $E$, and Cauchy–Schwarz on the RHS of (38) yields

$$
\int_0^T \left\| \mathbf{A}^{\frac{1}{2}} v_t \right\|_{L^2(\Gamma_1)}^2 dt \leq C \left[ \int_0^T \left( \left\| N^* A z_t \right\|_{L^2(\Gamma_1)}^2 + \left\| \mathbf{A}^{\frac{1}{2}} v_t \right\|_{L^2(\Gamma_1)}^2 \right) dt \right] \\
+ E(\mathbf{z}^\ast, \mathbf{v}^\ast, T) + E(\mathbf{z}^\ast, \mathbf{v}^\ast, 0) + \left\| v_t \right\|_{L^2(\Sigma_{1T})}^2,
$$

(39)

where $C$ is independent of time. The result follows upon adding (37) and (39), and recalling the characterization (10). □

**Lemma 2** There exists a constant $C$, independent of time, such that

$$
\int_{Q_T} z_t^2 \, dQ_T \leq C \left[ \int_0^T \left( \left\| z_t \right\|_{L^2(\Gamma_1)}^2 + \left\| \mathbf{A}^{\frac{1}{2}} v_t \right\|_{L^2(\Gamma_1)}^2 \right) dt \right] \\
+ \int_{\Sigma_{1T}} \left( \frac{\partial z}{\partial T} \right)^2 d\Sigma_{1T} + E(\mathbf{z}^\ast, \mathbf{v}^\ast, T) + E(\mathbf{z}^\ast, \mathbf{v}^\ast, 0) + \left\| z \right\|_{L^2(Q_T)}^2.
$$

(40)

**Proof.** With the given vector field $h(x)$ satisfying (h.i) and (h.iii), we have upon multiplying the wave equation in (1) by $h \cdot \nabla z$ the standard identity (see [14], Appendix A):

$$
\int_{Q_T} H \nabla z \cdot \nabla z \, dQ_T = \int_{\Sigma_T} \frac{\partial z}{\partial \nu} \cdot h \cdot \nabla z \, d\Sigma_T \\
+ \frac{1}{2} \int_{\Sigma_T} z_t^2 \, h \cdot \nu \, d\Sigma_T - \frac{1}{2} \int_{\Sigma_T} |\nabla z|^2 \, h \cdot \nu \, d\Sigma_T \\
- \frac{1}{2} \int_{Q_T} \left\{ z_t^2 - |\nabla z|^2 \right\} \, \text{div} \, h \, dQ_T - \left[ (z_t, h \cdot \nabla z)_{L^2(\Omega)} \right]^T_0.
$$

(41)

As $[z, z_t] \in D(\mathbf{A}^{\frac{1}{2}}) \times D(\mathbf{A}^{\frac{1}{2}})$, we then note that

on $\Sigma_{0T}$ $z_t = 0; \quad \frac{\partial z}{\partial \nu} = |\nabla z|; \quad h \cdot \nabla z = h \cdot \nu \frac{\partial z}{\partial \nu};$

and thus

$$
\int_{\Sigma_{0T}} \frac{\partial z}{\partial \nu} \cdot h \cdot \nabla z \, d\Sigma_{0T} + \frac{1}{2} \int_{\Sigma_{0T}} z_t^2 \, h \cdot \nu \, d\Sigma_{0T} - \frac{1}{2} \int_{\Sigma_{0T}} |\nabla z|^2 \, h \cdot \nu \, d\Sigma_{0T} \\
= \frac{1}{2} \int_{\Sigma_{0T}} |\nabla z|^2 \, h \cdot \nu \, d\Sigma_{0T} \leq 0,
$$

(42)

after using the condition (h.i). Inserting the inequality (42) into (41) will therefore yield

$$
\int_{Q_T} H \nabla z \cdot \nabla z \, dQ_T \leq \int_{\Sigma_{1T}} \frac{\partial z}{\partial \nu} \cdot h \cdot \nabla z \, d\Sigma_{1T} + \frac{1}{2} \int_{\Sigma_{1T}} z_t^2 \, h \cdot \nu \, d\Sigma_{1T} \\
- \frac{1}{2} \int_{\Sigma_{1T}} |\nabla z|^2 \, h \cdot \nu \, d\Sigma_{1T} - \frac{1}{2} \int_{Q_T} \left\{ z_t^2 - |\nabla z|^2 \right\} \, \text{div} \, h \, dQ_T \\
- \left[ (z_t, h \cdot \nabla z)_{L^2(\Omega)} \right]^T_0;
$$

(43)
hence, using the condition (h.iii), the Neumann B.C. in (1), the definition of $E$ and Cauchy–Schwarz gives us after estimating both sides of (43)

$$\rho \int_{Q_T} |\nabla z|^2 \, dQ_T \leq C \left[ \int_{\Sigma_{1T}} \left[ z_t^2 + v_t^2 \right] \, d\Sigma_{1T} + \int_{\Sigma_{1T}} |\nabla z|^2 \, d\Sigma_{1T} \right] + E(\overline{z}, \overline{v}, T) + E(\overline{z}, \overline{v}, 0) - \frac{1}{2} \int_{Q_T} \left[ z_t^2 - |\nabla z|^2 \right] \, \text{div} \, h \, dQ_T. \tag{44}$$

Now, to handle the last term on the RHS of (44), we multiply the wave equation (1) by $z \, \text{div} \, h$ and integrate by parts to obtain

$$\int_{Q_T} \left\{ z_t^2 - |\nabla z|^2 \right\} \, \text{div} \, h \, dQ_T = \left[ (z_t, z \, \text{div} \, h)_{L^2(\Omega)} \right]_0^T \tag{45}$$

$$+ \int_{Q_T} z \nabla (\text{div} \, h) \cdot \nabla z \, dQ_T - \int_{\Sigma_{1T}} \frac{\partial z}{\partial \nu} z \, \text{div} \, h \, d\Sigma_{1T},$$

after using Green’s Theorem and the identity $\nabla (z \, \text{div} \, h) \cdot \nabla z = z \nabla (\text{div} \, h) \cdot \nabla z + |\nabla z|^2 \, \text{div} \, h$. We thus have upon majorizing the RHS of (45) with the use of Poincaré’s inequality and the Neumann B.C. in (1),

$$\left| \int_{Q_T} \left\{ z_t^2 - |\nabla z|^2 \right\} \, \text{div} \, h \, dQ_T \right| \leq C_1 \left[ \int_{Q_T} z_t^2 \, dQ_T + \int_{\Sigma_{1T}} \left[ z_t^2 + v_t^2 \right] \, d\Sigma_{1T} \right] + E(\overline{z}, \overline{v}, T) + E(\overline{z}, \overline{v}, 0) + C_0 \int_{Q_T} |\nabla z|^2 \, dQ_T, \tag{46}$$

where $\epsilon > 0$ is arbitrarily small, and where the noncrucial dependence of $C_1$ upon $\epsilon$ has not been noted. Thus for $\epsilon$ small enough, adding the inequalities (44) and (46) together yields

$$(\rho - \epsilon C_0) \int_{Q_T} |\nabla z|^2 \, dQ_T \leq C_1 \left[ \int_{\Sigma_{1T}} \left[ z_t^2 + v_t^2 \right] \, d\Sigma_{1T} + \int_{\Sigma_{1T}} |\nabla z|^2 \, d\Sigma_{1T} \right]$$

$$+ E(\overline{z}, \overline{v}, T) + E(\overline{z}, \overline{v}, 0) + C_0 \int_{Q_T} z_t^2 \, dQ_T. \tag{47}$$

Moreover, (46) and (47) together give

$$(\rho - \epsilon C_0) \int_{Q_T} z_t^2 \, dQ_T \leq C_1 \left[ \int_{\Sigma_{1T}} \left[ z_t^2 + v_t^2 \right] \, d\Sigma_{1T} + \int_{\Sigma_{1T}} |\nabla z|^2 \, d\Sigma_{1T} \right]$$

$$+ E(\overline{z}, \overline{v}, T) + E(\overline{z}, \overline{v}, 0) + \int_{Q_T} z_t^2 \, dQ_T, \tag{48}$$

(where the constants $C_0$ and $C_1$ above are not necessarily the same throughout). Using (37) and the fact that $|\nabla z|^2 = \left( \frac{\partial z}{\partial \nu} \right)^2 + \left( \frac{\partial z}{\partial T} \right)^2$, we obtain the desired estimate (40) upon the addition of (47) and (48), and the use of the Neumann B.C. \(\square\)

With estimates (36) and (40) in conjunction with the relation

$$\int_0^T \left\| A^{\frac{1}{2}} z \right\|_{L^2(\Omega)}^2 \, dt = \int_{Q_T} z_t^2 - (z_t, z)_{L^2(\Omega)} \bigg|_0^T$$

$$- \int_0^T (z_t, z)_{L^2(\Gamma_1)} \, dt + \int_0^T (v_t, z)_{L^2(\Gamma_1)} \, dt \tag{49}$$

(obtained by multiplying (23) by $z$ and integrating), we then deduce the preliminary inequality

$$9$$
\[
\int_{\gamma}^{T} E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, t) \, dt \leq C \left[ \int_{0}^{T} \left[ \| z_t |_{\Gamma_1} \|_{L^2(\Gamma_1)}^2 + \| \overline{A}^\frac{1}{2} v_t \|_{L^2(\Gamma_1)}^2 \right] \, dt \right.
\]
\[+ \int_{\Sigma_{tr}} \frac{\partial z}{\partial t} \, d\Sigma_{tr} + E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, T) + E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, 0) + \| v \|_{L^2(\Sigma_{tr})}^2 + \| z \|_{L^2(Q_T)}^2 \right] ; \tag{50} \]

Repeating the same argument on the interval \((\epsilon, T - \epsilon)\), using the estimate (28) of Lemma A for the tangential derivative as well as the Neumann B.C. in (1), we arrive at
\[
\int_{\epsilon}^{T-\epsilon} E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, t) \, dt \leq C_T \left[ \int_{0}^{T} \left[ \| z_t |_{\Gamma_1} \|_{L^2(\Gamma_1)}^2 + \| \overline{A}^\frac{1}{2} v_t \|_{L^2(\Gamma_1)}^2 \right] \, dt \right.
\]
\[+ \| v \|_{L^2(\Sigma_{tr})}^2 + \| z \|_{H^\frac{1}{2}+ (Q_T)}^2 \right] + C \left[ E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, T) + E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, 0) \right], \tag{51} \]

where the constant \(C_T\) in (51) depends upon \(T\), but \(C\) does not. Using the relation (31) and its inherent dissipativity property, viz. \(E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, T) \leq E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, t) \forall \ 0 \leq t \leq T\), we have for \(T > 2C + 2\epsilon\),
\[
E(\overline{z}^\epsilon, \overline{\nu}^\epsilon, T) \leq \frac{C_T + 2C}{T - 2C - 2\epsilon} \left[ \int_{0}^{T} \left[ \| z_t |_{\Gamma_1} \|_{L^2(\Gamma_1)}^2 + \| \overline{A}^\frac{1}{2} v_t \|_{L^2(\Gamma_1)}^2 \right] \, dt \right.
\]
\[+ \| v \|_{L^2(\Sigma_{tr})}^2 + \| z \|_{H^\frac{1}{2}+ (Q_T)}^2 \right] . \tag{52} \]

So with (52) in hand, the proof of Theorem 1.1 will be complete if we can “absorb” the lower order terms \(\| v \|_{L^2(\Sigma_{tr})}^2\) and \(\| z \|_{H^\frac{1}{2}+ (Q_T)}^2\), which we now proceed to do.

**Lemma 3** Again, with the initial data \([\overline{z}_0^\epsilon, \overline{\nu}_0^\epsilon]\) in \(D(A)\) and with \(T\) sufficiently large, inequality (52) implies that there exists a positive constant \(C_T\) such that
\[
\| v \|_{C([0,T];L^2(\Gamma_1))} + \| z \|_{H^\frac{1}{2}+ (Q_T)} \leq C_T \left\{ \int_{0}^{T} \left[ \| z_t \|_{L^2(\Gamma_1)}^2 + \| \overline{A}^\frac{1}{2} v_t \|_{L^2(\Gamma_1)}^2 \right] \, dt \right\}. \tag{53} \]

**Proof:** We make use here of a compactness/uniqueness argument. If the lemma is false, then there exists a sequence \(\left\{ \overline{z}_n, \overline{v}_n \right\}_{n=1}^{\infty} \subseteq D(A)\), and a corresponding solution sequence \(\left\{ \overline{z}_n, \overline{v}_n \right\}_{n=1}^{\infty}\) which satisfies
\[
\| v_n \|_{C([0,T];L^2(\Gamma_1))} + \| z_n \|_{H^\frac{1}{2}+ (Q_T)} = 1 \forall \ n, \tag{54} \]
\[
\int_{0}^{T} \left[ \| z_n \|_{L^2(\Gamma_1)}^2 + \| \overline{A}^\frac{1}{2} v_n \|_{L^2(\Gamma_1)}^2 \right] \, dt \to 0 \text{ as } n \to \infty. \tag{55} \]

(55) and (52) then implies that the sequence \(\left\{ E(\overline{z}_n, \overline{v}_n, T) \right\}_{n=1}^{\infty}\) is bounded (uniformly in \(n\)), and consequently, (31) will have that \(\left\{ E(\overline{z}_n, \overline{v}_n, 0) \right\}_{n=1}^{\infty}\) is bounded. There then exists a subsequence, still denoted by \(\left\{ \overline{z}_n, \overline{v}_n \right\}_{n=1}^{\infty}\) and \(\overline{z}_0, \overline{v}_0 \in H_1 \times H_0\), such that
\[
\overline{z}_n \to \overline{z}_0 \text{ in } H_1 \text{ weakly;} \tag{56} \]
\[
\overline{v}_n \to \overline{v}_0 \text{ in } H_0 \text{ weakly.} \tag{57} \]
If we denote \([\tilde{z}, \tilde{v}, \tilde{u}, \tilde{v}_t] \equiv \begin{bmatrix} \tilde{z} \\ \tilde{v} \end{bmatrix}\) as the solution pair corresponding to the weak limits \([\tilde{z}_0, \tilde{v}_0]\), then \textit{a fortiori}

\[
\begin{bmatrix} z^{(n)}(t) \\ v^{(n)}(t) \end{bmatrix} \rightharpoonup \begin{bmatrix} \tilde{z} \\ \tilde{v} \end{bmatrix} \text{ in } L^\infty(0, T; H_1 \times H_0) \text{ weak star.}
\] (58)

Thus, \(z^{(n)} \rightharpoonup \tilde{z}\) weakly in \(H^1(Q_T)\), and consequently by the classic compactness theorem of Aubin.

\[
z^{(n)} \rightarrow \tilde{z} \text{ in } H^{1+\epsilon}(Q_T) \text{ strongly.}
\] (59)

Moreover, we deduce from (58) and a recent compactness result of Simon's (see [13], Corollary 4) that

\[
v^{(n)} \rightarrow \tilde{v} \text{ in } C([0, T]; L^2(\Gamma_1)) \text{ strongly.}
\] (60)

Consequently, taking the limit in (54),

\[
||\tilde{v}||_{C([0, T]; L^2(\Gamma_1))}^2 + ||\tilde{z}||_{H^{1+\epsilon}(Q_T)}^2 = 1 \quad \forall n.
\] (61)

Furthermore, the continuity of the map defined in \textbf{Proposition 2.1} and the convergence in (56)-(57) provide that

\[
\begin{align*}
\frac{\partial z^{(n)}}{\partial t} |_{\Gamma_1} & \rightarrow \tilde{z}_t |_{\Gamma_1} \text{ weakly in } L^2(\Sigma_1 T); \\
\tilde{A}^{1/2}v^{(n)} & \rightarrow \tilde{A}^{1/2}\tilde{v}_t \text{ weakly in } L^2(\Sigma_1 T);
\end{align*}
\] (62)\hspace{1cm} (63)

this convergence above, considered with that in (55) and the uniqueness of weak limits, allows one to deduce that

\[
\begin{align*}
\tilde{z}_t |_{\Gamma_1} &= 0; \\
\tilde{v}_t &= 0.
\end{align*}
\] (64)\hspace{1cm} (65)

From (65) then \(\tilde{v} = \text{constant}\), and combining this with the B.C. in (2) we have

\[
\tilde{v} = 0.
\] (66)

In dealing with the term \(\tilde{z}\), we bring forth the representation of the Hilbert space adjoint \(A^*\) given in [2] by

\[
A^* = \begin{bmatrix}
0 & -\mathbf{I} & 0 & 0 \\
A & -\alpha AN N^* A & 0 & -AN \\
0 & 0 & 0 & -\mathbf{I} \\
0 & N^* A & \tilde{A} & -\tilde{A}
\end{bmatrix},
\] (67)

\(\text{with } D(A^*) = \left\{[z_1, z_2, v_1, v_2] \in D(A^{1/2}) \times D(A^{1/2}) \times D(\tilde{A}^{1/2}) \times D(\tilde{A}^{1/2}) \\
such that } z_1 - \alpha NN^* Az_2 - Nv_2 \in D(A) \text{ and } v_1 - v_2 \in D(\tilde{A}) \right\}.\)

As \(\begin{bmatrix} \tilde{z} \\ \tilde{v} \end{bmatrix}\) is hence a weak solution of (1) from (65)-(66) and \(\tilde{z}_t |_{\Gamma_1} = 0\), we deduce from Ball's Theorem (see [3]) and the structure of \(A^*\) the following equation which holds pointwise for all \(z \in D(A^{1/2})\):

\[
\frac{d}{dt} \langle \tilde{z}_t(t), z \rangle_{L^2(\Omega)} = - \langle A^{1/2} \tilde{z}(t), A^{1/2} z \rangle_{L^2(\Omega)};
\] (68)
viz. \( \tilde{z}_t = -A\tilde{z} \in \big[ D(A^{1/2}) \big]' \), and if we make the change of variable \( p = A^{-1}\tilde{z}_t \) then \( p \in D(A) \), \( p|_{\Gamma_1} = 0 \), and so \( p \) solves the following wave equation:

\[
\begin{cases}
p_{tt} = \Delta p & \text{on } Q_T \\
p|_{\Gamma} = 0 & \text{on } \Sigma_T \\
\frac{\partial p}{\partial n}|_{\Gamma_1} = 0 & \text{on } \Sigma_{1T};
\end{cases}
\]

(69)

for \( T \) sufficiently large, we will hence have by Holmgren's Uniqueness Theorem that \( p = 0 \) which implies that \( \tilde{z} = 0 \), which coupled with (66) contradicts the equality (61), and the lemma is proved. \( \square \)

With Lemma 2.4 in hand, the proof of Theorem 1.1 is now complete.

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