

Inverse Problem for Kinetic Equations

Weiran Sun
Simon Fraser University

Collaborator:
Qin Li (U. Wisconsin)

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Outline

- Basic Setting;
- Examples: radiative transfer equations (RTE) in 2D
 - Recovery of the absorption coefficient: linear and nonlinear
 - Applications of energy estimates and maximum principle
 - Recovery of the scattering coefficient
 - Application of the velocity averaging lemma
- Future Questions.

Basic Setting

- Consider the stationary RTE

$$\mathbf{v} \cdot \nabla_x f = -\sigma_a f + F_f(x).$$

- $(x, \mathbf{v}) \in \Omega \times \mathbb{S}^1$.
- $\Omega \subseteq \mathbb{R}^2$ bounded, convex, C^1 boundary.
- $0 \leq \sigma_a = \sigma_a(x) \in L^\infty(\Omega)$.
- F_f depends on the moment(s) of f , where a k^{th} moment of f is given by

$$\int_{\mathbb{V}} \mathbf{v}^\alpha f \, d\mathbf{v}, \quad |\alpha| = k.$$

- $\langle f \rangle = \int_{\mathbb{S}^1} f \, d\mathbf{v}$: the zeroth moment of f .

Two Examples

- Linear RTE:

$$\mathbf{v} \cdot \nabla_x f = -\sigma_a f + \sigma_s \langle f \rangle,$$

where $0 \leq \sigma_s \leq \sigma_a$.

- Nonlinear RTE¹:

$$\mathbf{v} \cdot \nabla_x I = -\sigma_a I + \sigma_a T^4,$$

$$\Delta_x T = \sigma_a T^4 - \sigma_a \langle I \rangle.$$

where I is the intensity and T is the temperature.

¹Chandrasekhar, Radiative Transfer, Oxford University Press, 1950.

Main Framework

- Two components of our method:
 - Isolate the leading order behaviour: singular decomposition and/or iteration .²
 - Show that the remainder terms are regular enough to be of higher-order: energy method, velocity averaging lemma etc.

²Choulli, Stefanov, Uhlmann, Bal, Monard, Langmore, Ren, ...

Part I: Recover the Absorption Coefficient

Notation

- $\Gamma_- = \{(x, v) \mid x \in \partial\Omega, n(x) \cdot v < 0\} \subseteq \partial\Omega \times \mathbb{S}^1$.
- $\Gamma_+ = \{(x, v) \mid x \in \partial\Omega, n(x) \cdot v > 0\} \subseteq \partial\Omega \times \mathbb{S}^1$.
- $\tau_{\pm}(x, v)$: forward/backward exit times from point x along the direction v .

Absorption Coefficient: General RTE

- Equation

$$\mathbf{v} \cdot \nabla_x f = -\sigma_a f + F_f(\mathbf{x}), \quad (0.1)$$

Theorem (Li-S., 2019)

Let Ω be a convex and bounded domain with a C^1 boundary. Suppose $\sigma_a \geq 0$ and $\sigma_a \in C(\bar{\Omega})$. Suppose there exists $p > 1$ such that for any given incoming data ϕ satisfying

$$\phi \in L^p(\Gamma_-), \quad \phi \geq 0,$$

equation (0.7) has a unique nonnegative solution with the bound

$$\|F_f\|_{L^p(\Omega \times \mathbb{S}^1)} \leq C_0 \|\phi\|_{L^p(\Gamma_-)},$$

where C_0 is independent of ϕ and f . Then with proper choices of the incoming data, the absorption coefficient σ_a can be uniquely recovered from the measurement of the outgoing data.

Application: Linear RTE

- Two conditions to verify
 - Forward problem with incoming data is well-posed;
 - The forcing F_f is isotropic and for some $p > 1$,

$$\|F_f\|_{L^p(\Omega \times \mathbb{S}^1)} \leq C_0 \|\phi\|_{L^p(\Gamma_-)} .$$

- Forcing in the linear RTE

$$F_f(x) = \sigma_s(x) \langle f \rangle(x) .$$

- Given $\sigma_s \in L^\infty$, we have

$$\|F_f\|_{L^p(\Omega)} \leq \|\sigma_s\|_{L^\infty} \|f\|_{L^p(dx dv)} \leq \|\sigma_s\|_{L^\infty} \|\phi\|_{L^p(\Gamma_-)} .$$

- Conclusion: σ_a can be uniquely recovered in the linear RTE.

Application: Nonlinear RTE

- Equation with boundary data

$$v \cdot \nabla_x I = -\sigma_a I + \sigma_a T^4, \quad I|_{\Gamma_-} = \phi(x, v), \quad (0.2)$$

$$\Delta_x T = \sigma_a T^4 - \sigma_a \langle I \rangle, \quad T|_{\partial\Omega} = 0. \quad (0.3)$$

- Forcing: $F_f = \sigma_a T^4$.

Theorem (Well-posedness and bound of forcing)

Suppose $\phi \in L^\infty(\Gamma_-)$ and $\phi \geq 0$ Then

(a) System (0.2)-(0.3) has a unique solution satisfying that $I \geq 0, T \geq 0$,

$$\|I\|_{L^\infty(\Omega \times \mathbb{S}^1)} \leq \|\phi\|_{L^\infty(\Gamma_-)}, \quad \|T\|_{L^\infty(\Omega)} \leq \|\phi\|_{L^\infty(\Gamma_-)}^{1/4}.$$

(b) There exists a constant c_0 independent of ϕ such that

$$\|\sigma_a T^4\|_{L^2(\Omega)} \leq c_0 \|\phi\|_{L^2(\Gamma_-)}.$$

Nonlinear RTE: Well-posedness

- Main tools
 - Schauder fixed-point
 - Maximum principle for elliptic and RTE
 - Monotonicity argument for elliptic equations.
- Solution set

$$\mathcal{D} = \{T \mid 0 \leq T \leq \|\phi\|_{L^\infty}^{1/4}\} \subseteq L^\infty(\Omega).$$

- Mapping: $\mathcal{F} : \mathcal{D} \rightarrow \mathcal{D}$
 - Given $H \in \mathcal{D}$, let I_H be the solution such that

$$v \cdot \nabla_x I_H = -\sigma_a I_H + \sigma_a H^4, \quad I_H|_{\Gamma_-} = \phi(x, v).$$

- Define $\mathcal{F}(H) = T$ where T is the solution to the equation

$$\Delta_x T = \sigma_a T^4 - \sigma_a \langle I_H \rangle, \quad T|_{\partial\Omega} = 0.$$

Nonlinear RTE: Well-posedness

- Maximum principle for I_H :

$$v \cdot \nabla_x I_H = -\sigma_a I_H + \underbrace{\sigma_a H^4}_{\geq 0}, \quad I_H|_{\Gamma_-} = \phi(x, v) \geq 0.$$

which gives that $I_H \geq 0$.

- Maximum principle for $I_H - \|\phi\|_{L^\infty}$:

$$v \cdot \nabla_x (I_H - \|\phi\|_{L^\infty}) = -\sigma_a (I_H - \|\phi\|_{L^\infty}) + \underbrace{\sigma_a (H^4 - \|\phi\|_{L^\infty})}_{\leq 0},$$
$$(I_H - \|\phi\|_{L^\infty})|_{\Gamma_-} = \underbrace{\phi(x, v) - \|\phi\|_{L^\infty}}_{\leq 0}.$$

which implies that $I_H \leq \|\phi\|_{L^\infty}$.

- Overall, we have

$$0 \leq I_H \leq \|\phi\|_{L^\infty}.$$

Nonlinear RTE: Well-posedness

- Equation for T (given $0 \leq I_H \leq \|\phi\|_{L^\infty}$)

$$-\Delta_x T = -\sigma_a T^4 + \sigma_a \langle I_H \rangle, \quad T|_{\partial\Omega} = 0. \quad (0.4)$$

- Monotonicity Method

- Find an increasing sequence T_k

$$0 \leq \underline{T} \leq T_1 \leq \dots \leq T_k \leq \dots \leq \bar{T}$$

where \underline{T} is a subsolution and \bar{T} is a supersolution.

- Show that $T_k \rightarrow T$ and T solves (0.4).
- Maximum principle: at x_0 where T attains its maximum, we have

$$0 \leq -\Delta_x T(x_0) = -\sigma_a T^4(x_0) + \sigma_a \langle I_H \rangle(x_0),$$

which implies $T \leq \|\phi\|_{L^\infty}^{1/4}$.

Bound of Forcing

- Recall the equation for (I, T)

$$v \cdot \nabla_x I = -\sigma_a I + \sigma_a T^4, \quad I|_{\Gamma_-} = \phi(x, v), \quad (0.5)$$

$$\Delta_x T = \sigma_a T^4 - \sigma_a \langle I \rangle, \quad T|_{\partial\Omega} = 0. \quad (0.6)$$

- We use the energy method to show that

$$\|\sigma_a T^4\|_{L^2(\Omega)} \leq c_0 \|\phi\|_{L^2(\Gamma_-)}.$$

- Basic energy estimate³

$$\frac{1}{2} \int_{\partial\Omega} \int_{\mathbb{S}^1} (n(x) \cdot v) I^2 + 4 \int_{\Omega} T^3 |\nabla_x T|^2 + \int_{\Omega} \int_{\mathbb{S}^1} \sigma_a (\langle I \rangle - T^4)^2 \leq 0.$$

which implies that

$$\begin{aligned} \int_{\Omega} \int_{\mathbb{S}^1} \sigma_a (\langle I \rangle - T^4)^2 \, dx \, dv &\leq -\frac{1}{2} \int_{\partial\Omega} \int_{\mathbb{S}^1} (n(x) \cdot v) I^2 \, dv \, dS_x \\ &= -\frac{1}{2} \left(\iint_{\Gamma_+} + \iint_{\Gamma_-} \right) \leq \frac{1}{2} \|\phi\|_{\Gamma_-}^2. \end{aligned}$$

³Klar-Schmeiser, M3AS, 2001.

Bound of Forcing

- Rewrite the I -equation as

$$v \cdot \nabla_x I = -\sigma_a(I - \langle I \rangle) + \underbrace{\sigma_a(T^4 - \langle I \rangle)}_{\text{forcing}}.$$

- Energy Estimate:

$$\|I\|_{L^2(\Omega \times \mathbb{S}^1)} \leq c_0 \left(\|\sigma_a(T^4 - \langle I \rangle)\|_{L^2(\Omega)} + \|\phi\|_{L^2(\Gamma_-)} \right) \leq c_0 \|\phi\|_{L^2(\Gamma_-)}$$

- Combine the above estimates

$$\begin{aligned} \|\sigma_a T^4\|_{L^2(\Omega)} &\leq \|\sigma_a(T^4 - \langle I \rangle)\|_{L^2(\Omega)} + \|\langle I \rangle\|_{L^2(\Omega)} \\ &\leq \|\sigma_a(T^4 - \langle I \rangle)\|_{L^2(\Omega)} + \|I\|_{L^2(\Omega \times \mathbb{S}^1)} \\ &\leq c_0 \|\phi\|_{L^2(\Gamma_-)}. \end{aligned}$$

which gives the desired forcing bound.

Absorption Coefficient: General RTE

- Equation

$$\mathbf{v} \cdot \nabla_x f = -\sigma_a f + F_f(\mathbf{x}), \quad (0.7)$$

Theorem (Li-S., 2019)

Let Ω be a convex and bounded domain with a C^1 boundary. Suppose $\sigma_a \geq 0$ and $\sigma_a \in C(\bar{\Omega})$. Suppose there exists $p > 1$ such that for any given incoming data ϕ satisfying

$$\phi \in L^p(\Gamma_-), \quad \phi \geq 0,$$

equation (0.7) has a unique nonnegative solution with the bound

$$\|F_f\|_{L^p(\Omega \times \mathbb{S}^1)} \leq C_0 \|\phi\|_{L^p(\Gamma_-)},$$

where C_0 is independent of ϕ and f . Then with proper choices of the incoming data, the absorption coefficient σ_a can be uniquely recovered from the measurement of the outgoing data.

Idea of Proof

- Incoming data ϕ :

$$\phi(x, v) = \frac{1}{\epsilon\delta} \phi_0\left(\frac{|x - x^{in}|}{\epsilon}\right) \phi_0\left(\frac{|v - v^{in}|}{\delta}\right), \quad (x, v) \in \Gamma_-.$$

where

- $\phi_0 \in C^\infty(\mathbb{R})$ is a compactly supported bump function

$$\phi_0(r) = \begin{cases} 1, & |r| \leq 1, \\ 0, & |r| \geq 2, \\ \text{smooth on } \mathbb{R}. \end{cases}$$

- $\int_{\mathbb{R}} \phi_0(r) dr = 1.$
- $(x^{in}, v^{in}) \in \Gamma_-$ such that $n(x^{in}) \cdot v^{in} < 0.$
- Bound: $\|\phi\|_{L^p(\Gamma_-)} \leq c_0 \epsilon^{-(1-1/p)} \delta^{-(1-1/p)}$ for any $p \geq 1.$

Measurement

- Total measurement

$$M_\psi = \iint_{\Gamma_+} \psi(x, v) f(x, v) d\Gamma_+, \quad d\Gamma = n(x) \cdot v dS_x dv.$$

where

- the measurement function ψ is

$$\psi(x, v) = \psi_0(x - x^{out}) \psi_0\left(\frac{v - v^{out}}{\delta}\right) \quad (x, v) \in \Gamma_+.$$

- Choice of (x^{out}, v^{out}) :

$$x^{out} = x^{in} + \tau_+(x^{in}, v^{in}) v^{in}, \quad v^{out} = v^{in}.$$

Decomposition

- Full equation

$$\begin{aligned}v \cdot \nabla_x f &= -\sigma_a f + F_f(x), & x \in \Omega \subseteq \mathbb{R}^2, \\ f|_{\Gamma_-} &= \phi(x, v), & (x, v) \in \Gamma_-.\end{aligned}$$

- Decompose the solution as $f = f_1 + f_2$.
- Leading-order

$$\begin{aligned}v \cdot \nabla_x f_1 &= -\sigma_a f_1, & x \in \Omega \subseteq \mathbb{R}^2, \\ f_1|_{\Gamma_-} &= \phi(x, v), & (x, v) \in \Gamma_-.\end{aligned}$$

- Remainder

$$\begin{aligned}v \cdot \nabla_x f_2 &= -\sigma_a f_2 + F_f(x), & x \in \Omega \subseteq \mathbb{R}^2, \\ f_2|_{\Gamma_-} &= 0, & (x, v) \in \Gamma_-.\end{aligned}$$

Leading Order

- Formula for f_1 :

$$f_1(x, v) = e^{-\int_0^{\tau_-(x,v)} \sigma_a(x-sv) ds} \phi(x - \tau_-(x, v)v, v),$$

for all $(x, v) \in \bar{\Omega} \times \mathbb{S}^1$.

- Contribution of f_1 to the measurement:

$$M_\psi(f_1) = \iint_{\Gamma_+} \psi(x, v) f_1(x, v) d\Gamma_+ = \frac{1}{\epsilon} \int_{\partial\Omega} \psi_0(x - x^{out}) G_{\epsilon, \delta}(x) dS_x.$$

where

$$G_{\epsilon, \delta}(x) = \frac{1}{\delta} \int_{\mathbb{S}_{x,+}^1} e^{-\int_0^{\tau_-(x,v)} \sigma_a(x-sv) ds} \psi_0\left(\frac{|v - v^{in}|}{\delta}\right) \\ \times \phi_0\left(\frac{|x - \tau_-(x, v)v - x^{in}|}{\epsilon}\right) \phi_0\left(\frac{|v - v^{in}|}{\delta}\right) n(x) \cdot v dv.$$

- Contribution from f_1 :

$$\lim_{\epsilon \rightarrow 0} \lim_{\delta \rightarrow 0} M_\psi(f_1) = C_{\psi_0, \phi_0} e^{-\int_0^{\tau_-(x^{out}, v^{out})} \sigma_a(x^{out} - sv^{out}) ds}.$$

Remainder

- Goal: show that $\lim_{\epsilon \rightarrow 0} \lim_{\delta \rightarrow 0} M_\psi(f_2) = 0$.
- Formula for f_2

$$f_2(x, v) = \int_0^{\tau_-(x, v)} e^{-\int_0^s \sigma_a(x - \tau v) d\tau} F_f(x - sv) ds,$$

for all $(x, v) \in \bar{\Omega} \times \mathbb{S}^1$.

- Contribution from f_2

$$\begin{aligned} M_\psi(f_2) &= \iint_{\Gamma_+} \psi(x, v) f_2(x, v) d\Gamma_+ \\ &= \iint_{\Gamma_+} \int_0^{\tau_-(x, v)} \psi(x, v) e^{-\int_0^s \sigma_a(x - \tau v) d\tau} F_f(x - sv) n(x) \cdot v ds dS_x dv, \end{aligned}$$

- Main ingredients for smallness of $M_\psi(f_2)$
 - Regularity and isotropy of F_f
 - Smallness of the support of integration.

Remainder

- Recall that by assumption,

$$\|F_f\|_{L^p(\Omega \times \mathbb{S}^1)} \leq C_0 \|\phi\|_{L^p(\Gamma_-)}, \quad \psi(x, v) = \psi_0(x - x^{out}) \psi_0\left(\frac{v - v^{out}}{\delta}\right).$$

- Change of variables: $y = x - sv$ such that

$$\iint_{\Gamma_+} \int_0^{\tau^-(x,v)} \dots n(x) \cdot v \, ds \, dS_x \, dv = \int_{\mathbb{S}^1} \int_{\Omega} \dots \, dy \, dv.$$

- Estimate

$$\begin{aligned} & \left| \iint_{\Gamma_+} \psi(x, v) f_2(x, v) \, d\Gamma \right| \\ &= \iint_{\Gamma_+} \int_0^{\tau^-(x,v)} \psi(x, v) e^{-\int_0^s \sigma_a(x - \tau v) \, d\tau} F_f(x - sv) n(x) \cdot v \, ds \, dS_x \, dv \\ &\leq \int_{\mathbb{S}^1} \int_{\Omega} \psi_0\left(\frac{v - v^{in}}{\delta}\right) |F_f(y)| \, dy \, dv && \text{(Change of Variable)} \\ &= \left(\int_{\mathbb{S}^1} \psi_0\left(\frac{v - v^{in}}{\delta}\right) \, dv \right) \left(\int_{\Omega} |F_f(y)| \, dy \right) && \text{(Isotropy)} \\ &\leq c_0 \delta \|F_f\|_{L^p(\Omega)} \leq c_0 \delta \|\phi\|_{L^p(\Gamma_-)} \leq c_0 \epsilon^{-(1-1/p)} \delta^{1/p}. && \text{(Regularity)} \end{aligned}$$

- $\lim_{\epsilon \rightarrow 0} \lim_{\delta \rightarrow 0} M_\psi(f_2) = 0.$

Part II: Recover the Scattering Coefficient in 2D

RTE: Scattering Coefficient

- Equation

$$\mathbf{v} \cdot \nabla_x f = -\sigma_a f + \sigma_s \langle f \rangle, \quad (0.8)$$

Theorem (Li-S., 2019)

Let $\Omega \subseteq \mathbb{R}^2$ be a convex and bounded domain with a C^1 boundary. Suppose $\sigma_a \in C(\bar{\Omega})$ is given and

$$\sigma_s = \sigma_s(x), \quad 0 < \sigma_0 \leq \sigma_s \leq \sigma_a.$$

Then with proper choices of the incoming data, the scattering coefficient σ_s in (0.8) can be uniquely recovered from the measurement of the outgoing data.

Main Decomposition

- Decompose the solution f into three parts: $f = f_1 + f_2 + f_3$ where

$$\mathbf{v} \cdot \nabla_x f_1 = -\sigma_a f_1, \quad f_1|_{\Gamma_-} = \phi(x, \mathbf{v}),$$

$$\mathbf{v} \cdot \nabla_x f_2 = -\sigma_a f_2 + \sigma_s \langle f_1 \rangle, \quad f_2|_{\Gamma_-} = 0,$$

$$\mathbf{v} \cdot \nabla_x f_3 = -\sigma_a f_3 + \sigma_s \langle f_3 \rangle + \sigma_s \langle f_2 \rangle, \quad f_3|_{\Gamma_-} = 0,$$

- Let M_ψ be the measurement of the outgoing data.
- Main idea: Choose ϕ and ψ in the way such that
 - $M_\psi(f_1) = 0$;
 - $M_\psi(f_2) = \mathcal{O}(1)$;
 - $M_\psi(f_3)$ is of higher order due to the regularity of f_3 .
- Use the averaging lemma to show that $\langle f_2 \rangle$ is more regular than f_2 .

Averaging Lemma (Simplest Case)

Theorem (Golse-Lions-Perthame-Sentis, JFA, 1988)

Suppose $0 < \sigma_0 \leq \sigma_s \leq \sigma_a$ with $\sigma_a \in C(\bar{\Omega})$. Suppose $\phi \in L^p(\Gamma_-)$ for some $p > 1$ and f satisfies the equation

$$v \cdot \nabla_x f = -\sigma_a f + \sigma_s \langle f \rangle + g, \quad f|_{\Gamma_-} = \phi(x, v). \quad (0.9)$$

Then for any $\gamma \leq \inf\{\frac{1}{p}, 1 - \frac{1}{p}\}$, the velocity average of f satisfies $\langle f \rangle \in W^{\gamma,p}(\Omega)$ with the bound

$$\|\langle f \rangle\|_{W^{\gamma,p}(\Omega)} \leq C_0 \left(\|\phi\|_{L^p(\Gamma_-)} + \|g\|_{L^p(\Omega \times S^1)} \right).$$

In particular, by Sobolev embedding, we have $\langle f \rangle \in L^q(\Omega)$ where

$$q = \frac{1}{\frac{1}{p} - \frac{\gamma}{d}} > p.$$

Application of Averaging Lemma

- Equations for f_1, f_2, f_3 :

$$\begin{aligned}v \cdot \nabla_x f_1 &= -\sigma_a f_1, & f_1|_{\Gamma_-} &= \phi(x, v), \\v \cdot \nabla_x f_2 &= -\sigma_a f_2 + \sigma_s \langle f_1 \rangle, & f_2|_{\Gamma_-} &= 0, \\v \cdot \nabla_x f_3 &= -\sigma_a f_3 + \sigma_s \langle f_3 \rangle + \sigma_s \langle f_2 \rangle, & f_3|_{\Gamma_-} &= 0,\end{aligned}$$

- Estimates using the averaging lemma:

$$\begin{aligned}\| \langle f_1 \rangle \|_{L^{p_1}(\Omega)} &\leq c_0 \| \langle f_1 \rangle \|_{W^{s_0, r}} \leq c_0 \| \phi \|_{L^r(\Gamma_-)} \leq c_0 \epsilon^{-\frac{r-1}{r}} \delta^{-\frac{r-1}{r}} \eta^{\frac{1}{r}}, \\ \| \langle f_2 \rangle \|_{L^{p_2}(\Omega)} &\leq c_0 \| \langle f_2 \rangle \|_{W^{s_1, p_1}(\Omega)} \leq c_0 \| \langle f_1 \rangle \|_{L^{p_1}(\Omega \times \mathbb{S}^1)} \leq c_0 \epsilon^{-\frac{r-1}{r}} \delta^{-\frac{r-1}{r}} \eta^{\frac{1}{r}},\end{aligned}$$

where

$$\begin{aligned}s_0 &= \inf \left\{ \frac{1}{r}, 1 - \frac{1}{r} \right\}, & s_1 &= \inf \left\{ \frac{1}{p_1}, 1 - \frac{1}{p_1} \right\}, \\ p_1 &= \frac{1}{\frac{1}{r} - \frac{s_0}{2}} > r, & p_2 &= \frac{1}{\frac{1}{p_1} - \frac{s_1}{2}} > p_1 > r.\end{aligned}$$

Data and Measurement

- Let $(x^{in}, v^{in}) \in \Gamma_-$ and $(x^{out}, v^{out}) \in \Gamma_+$ such that $v^{in} \nparallel v^{out}$.
- Let ϕ and ψ be given by ⁴

$$\phi(x, v) = \frac{1}{\epsilon\delta} \phi_0\left(\frac{(x - x^{in}) \cdot v_{\perp}^{in}}{\epsilon\eta}\right) \phi_0\left(\frac{|v - v^{in}|}{\delta}\right), \quad (x, v) \in \Gamma_-,$$
$$\psi(x, v) = \frac{1}{\theta\beta} \psi_0\left(\frac{|x - x^{out}|}{\theta}\right) \psi_0\left(\frac{|v - v^{out}|}{\beta}\right), \quad (x, v) \in \Gamma_+.$$

where ϕ_0, ψ_0 are bump functions and

$$v_{\perp}^{in} \cdot v^{in} = 0 \quad \text{and} \quad \eta = v_{\perp}^{in} \cdot v^{out} > 0.$$

- For $r > 1$ to be chosen later, we have

$$\|\phi\|_{L^r(\Gamma_-)} \leq c_0 \epsilon^{-\frac{r-1}{r}} \delta^{-\frac{r-1}{r}} \eta^{\frac{1}{r}}, \quad r > 1.$$

⁴Choulli-Stefanov, Osaka J. Math., 1998.

Contribution from f_3

- Equation

$$v \cdot \nabla_x f_3 = -\sigma_a f_3 + \underbrace{\sigma_s \langle f_3 \rangle + \sigma_s \langle f_2 \rangle}_{\text{forcing}}, \quad f_3|_{\Gamma_-} = 0,$$

- Contribution to the measurement $M_\psi(f_3)$

$$\begin{aligned} \left| \iint_{\Gamma_+} \psi(x, v) f_3(x, v) d\Gamma \right| &\leq c_0 \int_{\mathbb{S}^1} \int_{\Omega} \psi(y + \tau_+(y, v)v) |\langle f_2 \rangle(y) + \langle f_3 \rangle(y)| dy dv \\ &\leq \frac{c_0}{\theta} \left(\int_{\partial\Omega} \psi_0^{p'_2} \left(\frac{|x - x^{out}|}{\theta} \right) dS_x \right)^{1/p'_2} \|\langle f_2 \rangle\|_{L^{p_2}(\Omega)} \\ &\leq c_0 \theta^{-\frac{1}{p_2}} \epsilon^{-\frac{r-1}{r}} \delta^{-\frac{r-1}{r}} \eta^{\frac{1}{r}}. \end{aligned}$$

- We will choose proper parameters to make $M_\psi(f_3)$ small.

Contribution from f_1

- Equation

$$\mathbf{v} \cdot \nabla_x f_1 = -\sigma_a f_1, \quad f_1|_{\Gamma_-} = \phi(\mathbf{x}, \mathbf{v}),$$

- Contribution to the measurement $M_\psi(f_1)$

$$\begin{aligned} \mathcal{M}_1 = \frac{1}{\epsilon\delta} \frac{1}{\theta\beta} \iint_{\Gamma_+} e^{-\int_0^{\tau_-(x,v)} \sigma_a(x-sv) ds} \psi_0\left(\frac{|\mathbf{x} - \mathbf{x}^{out}|}{\theta}\right) \psi_0\left(\frac{|\mathbf{v} - \mathbf{v}^{out}|}{\beta}\right) \\ \times \phi_0\left(\frac{(\mathbf{x} - \tau_-(\mathbf{x}, \mathbf{v})\mathbf{v} - \mathbf{x}^{in}) \cdot \mathbf{v}_\perp^{in}}{\epsilon\eta}\right) \phi_0\left(\frac{|\mathbf{v} - \mathbf{v}^{in}|}{\delta}\right) d\Gamma_+. \end{aligned}$$

- Make $M_\psi(f_1) = 0$ by letting

$$\text{supp}\left(\psi_0\left(\frac{|\mathbf{x} - \mathbf{x}^{out}|}{\theta}\right)\right) \cap \text{supp}\left(\phi_0\left(\frac{(\mathbf{x} - \tau_-(\mathbf{x}, \mathbf{v})\mathbf{v} - \mathbf{x}^{in}) \cdot \mathbf{v}_\perp^{in}}{\epsilon\eta}\right)\right) = \emptyset.$$

- Sufficient condition

$$|\mathbf{x}^{out} - \mathbf{x}_0^{out}| \geq c_0(\theta + \delta + \epsilon\eta).$$

Contribution from f_2

- Equation

$$\mathbf{v} \cdot \nabla_x f_2 = -\sigma_a f_2 + \sigma_s \langle f_1 \rangle, \quad f_2|_{\Gamma_-} = 0,$$

- Contribution to the measurement $M_\psi(f_2)$

$$\begin{aligned} M_\psi(f_2) &= \iint_{\Gamma_+} \int_0^{\tau_-(x,v)} \psi(x,v) e^{-\int_0^s \sigma_a(x-\tau v) d\tau} \sigma_s(x-sv) \langle f_1 \rangle(x-sv) ds d\Gamma \\ &\sim \frac{1}{\epsilon\delta} \frac{1}{\theta\beta} \iint_{\Gamma_+} \int_0^{\tau_-(x,v)} \int_{\mathbb{S}^1} H(s, x^{out}, v^{out}, v^{in}) \psi_0 \left(\frac{|x-x^{out}|}{\theta} \right) \psi_0 \left(\frac{|v-v^{out}|}{\beta} \right) \\ &\quad \times \phi_0 \left(\frac{((x-sv)'_w - x^{in}) \cdot v_\perp^{in}}{\epsilon\eta} \right) \phi_0 \left(\frac{|w-v^{in}|}{\delta} \right) dw ds d\Gamma_+ \end{aligned}$$

- $H(s, x, v, w) = e^{-\int_0^s \sigma_a(x-\tau v) d\tau} e^{-\int_0^{\tau_-(x-sv,w)} \sigma_a(x-sv-\tau w) d\tau} \sigma_s(x-sv).$
- $(x-sv)'_v$ is the backward exit point of $x-sv$ along the direction v .

Contribution from f_2

- Focus on the main term

$$\phi_0 \left(\frac{((x - sv)'_w - x^{in}) \cdot v_{\perp}^{in}}{\epsilon \eta} \right), \quad \eta = v_{\perp}^{in} \cdot v^{out} > 0.$$

- Given that $(x, v, w) \sim (x^{out}, v^{out}, v^{in})$, we have

$$(x - sv)'_w \sim (x^{out} - sv^{out}) - \tau_{-}(x^{out} - sv^{out}, v^{in})v^{in}$$

Together with

$$x^{in} = (x^{out} - s_0 v^{out}) - \tau_{-}(x^{out} - s_0 v^{out}, v^{in})v^{in}$$

we have

$$\frac{|((x - sv)'_w - x^{in}) \cdot v_{\perp}^{in}|}{\epsilon \eta} = \frac{|s - s_0|}{\epsilon}.$$

- $\lim_{\epsilon, \delta, \theta, \beta \rightarrow 0} \mathcal{M}_{2,1} = e^{-\int_0^{\tau^-(x_0, v^{in})} \sigma_a(x_0 - \tau x^{in}) d\tau} \sigma_s(x_0).$

Choice of Parameters

- Constraints of parameters

$$\theta^{-\frac{1}{p_2}} \epsilon^{-\frac{r-1}{r}} \delta^{-\frac{r-1}{r}} \eta^{\frac{1}{r}} \ll 1, \quad \theta + \delta + \beta \ll \epsilon \eta, \quad \beta \ll \theta. \quad (0.10)$$

- Borderline case: let

$$\theta = \epsilon \eta, \quad \theta = \delta = \beta.$$

In this case we require

$$\frac{1}{p_2} + \frac{2(r-1)}{r} < 1.$$

- Example: take $r \in (1, 3/2)$ to obtain that

$$s_0 = 1 - \frac{1}{r}, \quad p_1 = \frac{2r}{3-r} < 2, \quad s_2 = 1 - \frac{1}{p_1}, \quad \frac{1}{p_2} = \frac{3-p_1}{2p_1} = \frac{9}{4r} - \frac{5}{4},$$

which gives

$$\frac{1}{p_2} + \frac{2(r-1)}{r} = 1 + \frac{1}{4} \left(\frac{1}{r} - 1 \right) < 1.$$

Future Questions

- To name a few...
 - Time-dependent transport equations;
 - Other types of kinetic equations: BGK, Vlasov-type.
 - Multi-scale problems.